Medical Equipment II - 2010

Chapter 15: Interaction of Photons and Charged Particles with Matter(2)

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Photoelectric Effect

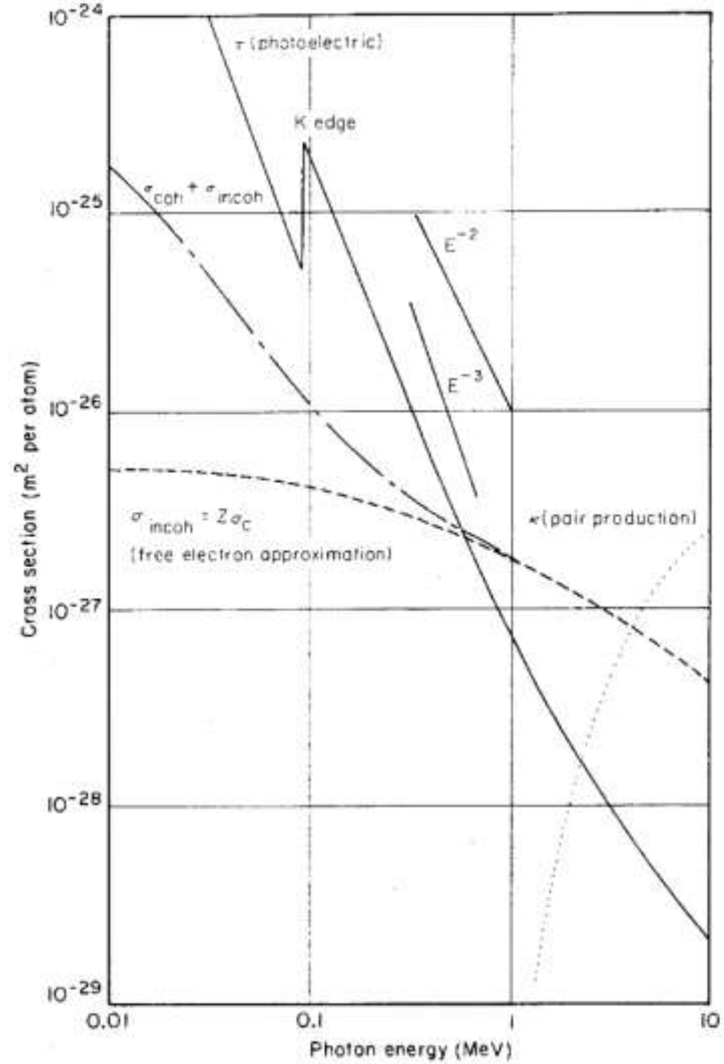
- (γ, *e*) Photon interaction – $h\nu_0 = T_{\mathrm{el}} + B$
  - T_{el}: Kinetic energy of electron, B: binding energy
- Binding energy depends on shell
  - B_K, B_L, and so on.
- Photoelectric cross section is τ .

$$\tau = \tau_K + \tau_L + \tau_M + \cdots .$$

Photoelectric Effect

- For photon energies too small to remove an electron from the *K* shell, τ_K is zero.
  - K edge
  - Can still remove L electron
- Model around 100 KeV:

$$\tau \propto Z^4 E^{-3}.$$



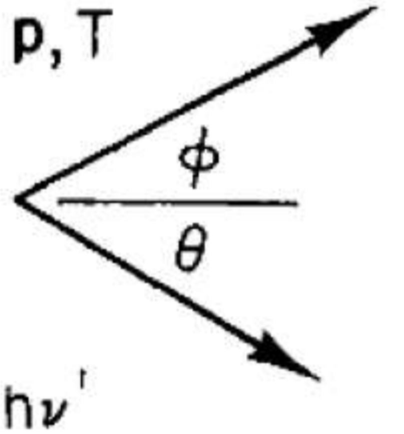
Compton Scattering: Kinematics

- (γ, γ' *e)* photon interaction $h\nu_0 = h\nu + T_{\mathrm{el}} + B.$
- Photon kinematics: Special relativity

$$E^2 = (pc)^2 + (m_0 c^2)^2. \quad \Rightarrow \quad E = h\nu = pc.$$

- Conservation of energy and momentum can be used to derive angle and energy of scattered photon

Compton Scattering: Kinematics



- Conservation of momentum in direction of the incident photon:

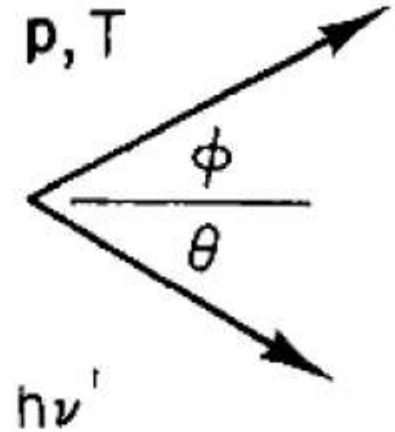
$$\frac{h\nu_0}{c} = \frac{h\nu'}{c}\cos\theta + p\cos\phi,$$

- Conservation of momentum at 90°

$$\frac{h\nu'}{c}\sin\theta = p\sin\phi.$$

- Conservation of energy $$h\nu_0 = h\nu' + T.$$

Compton Scattering: Kinematics



- Electron energy:

$$E = T + m_e c^2$$

- Combining with special relativity:

$$E^2 = (pc)^2 + (m_0 c^2)^2.$$

$$(pc)^2 = T^2 + 2 m_e c^2 T.$$

- Solve 4 equations in 4 unknowns
  - Unknowns: ***T*, *ν'*, *θ*, *ϕ***

Compton Scattering: Kinematics

- Wavelength of scattered photon:

$$\lambda' - \lambda_0 = \frac{c}{\nu'} - \frac{c}{\nu_0} = \frac{h}{m_e c}(1 - \cos\theta).$$

  - Difference is independent of incident wavelength
  - Compton length of electron $h/m_e c$
- Energy of scattered photon

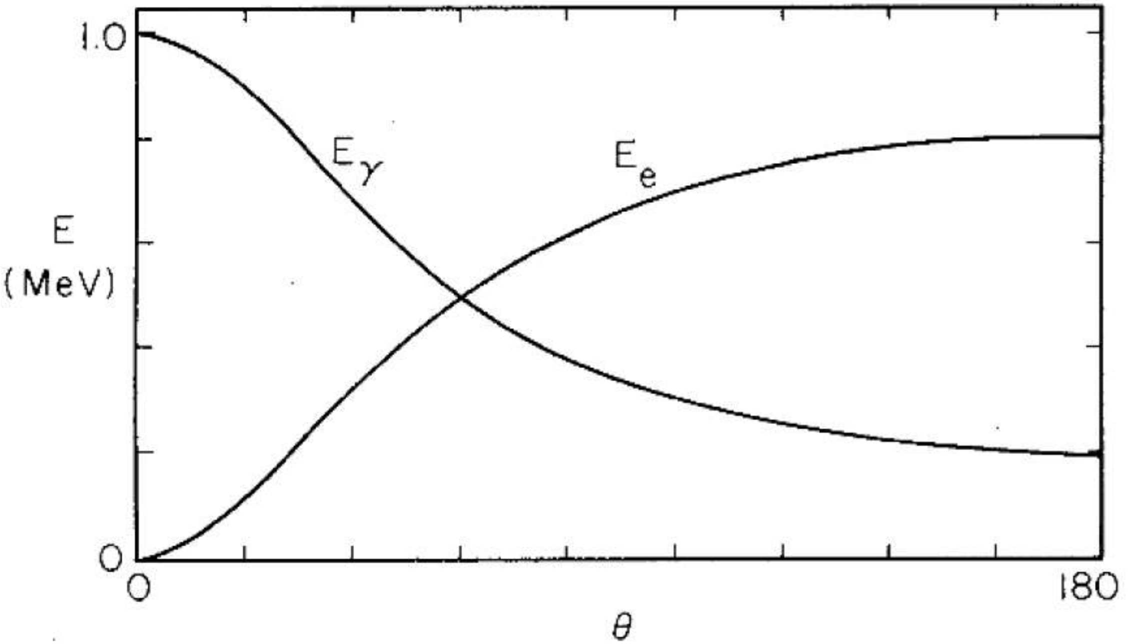
$$h\nu' = \frac{m_e c^2}{1 - \cos\theta + 1/x} \qquad x = \frac{h\nu_0}{m_e c^2}.$$

Compton Scattering: Kinematics

- Energy of recoil electron

$$T = \frac{h\nu_0(2x\cos^2\phi)}{(1+x)^2 - x^2\cos^2\phi} = \frac{h\nu_0 x(1-\cos\theta)}{1+x(1-\cos\theta)}.$$

- Dependence on angle θ



Compton Scattering: Cross Section

- Compton cross section σ_C.
- Quantum mechanics: Klein–Nishina Formula

$$\frac{d\sigma_C}{d\Omega} = \frac{r_e^2}{2}\left[\frac{1+\cos^2\theta + \dfrac{x^2(1-\cos\theta)^2}{1+x(1-\cos\theta)}}{[1+x(1-\cos\theta)]^2}\right]$$

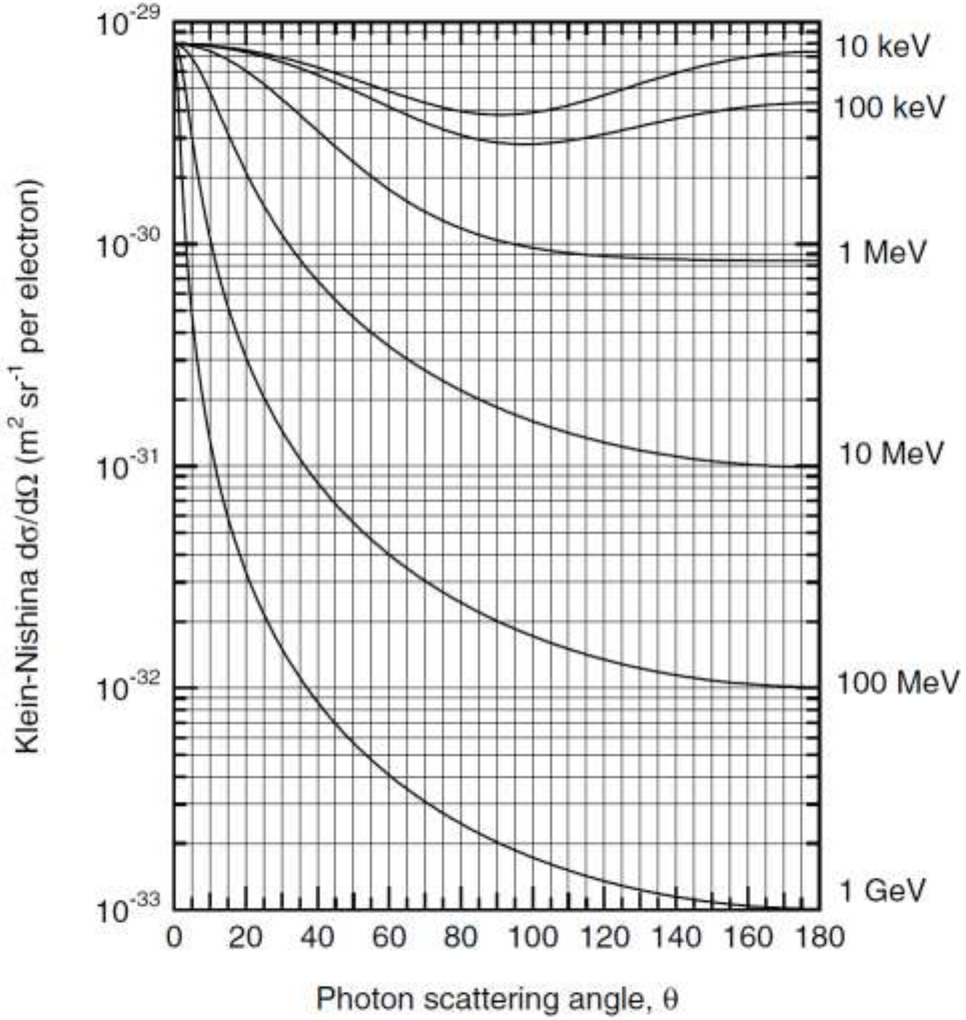
  - Classical radius of electron

$$r_e = \frac{e^2}{4\pi\epsilon_0 m_e c^2} = 2.818\times 10^{-15}\ \text{m},$$

Compton Scattering: Cross Section

- σ_C peaked in the forward direction at high energies.
- As $x \rightarrow 0$ (high energy):

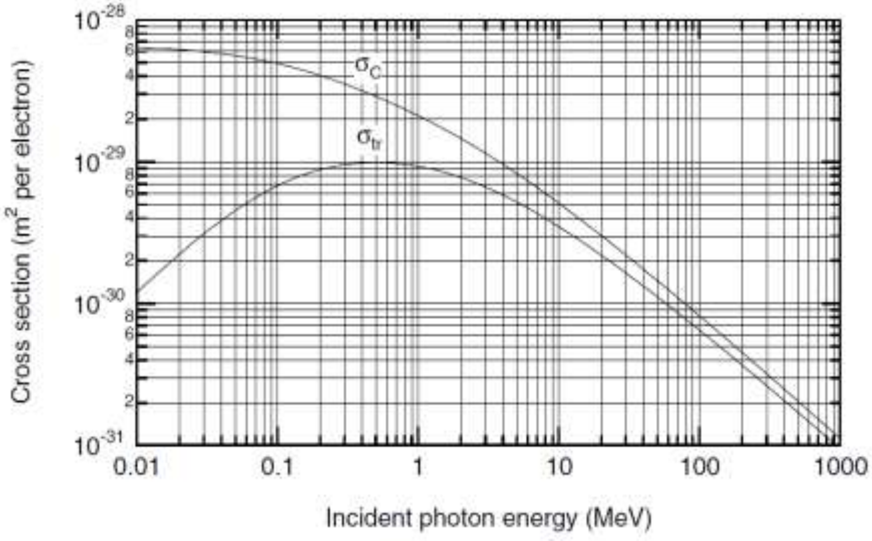
$$\frac{d\sigma_C}{d\Omega} = \frac{r_e^2(1+\cos^2\theta)}{2}$$



Compton Scattering: Cross Section

- Integrated over all angles

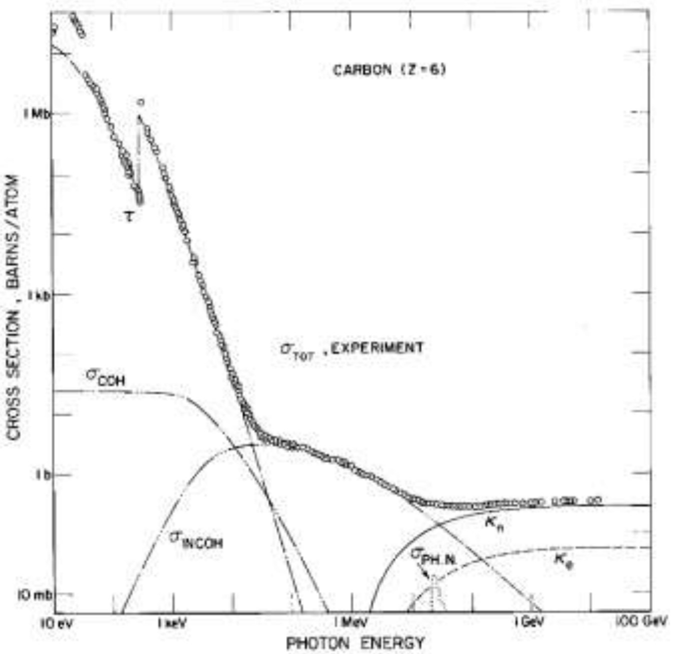
$$\sigma_C = 2\pi r_e^2 \left[\frac{1+x}{x^2} \left(\frac{2(1+x)}{1+2x} - \frac{\ln(1+2x)}{x} \right) + \frac{\ln(1+2x)}{2x} - \frac{1+3x}{(1+2x)^2} \right]. \quad (15.19)$$



Compton Scattering: Incoherent Scattering

- σ_C is for a single electron.
- For an atom containing Z electrons, maximum value of σ_{incoh} occurs if all Z electrons take part in Compton scattering

$$\sigma_{\text{incoh}} \leq Z\sigma_C.$$



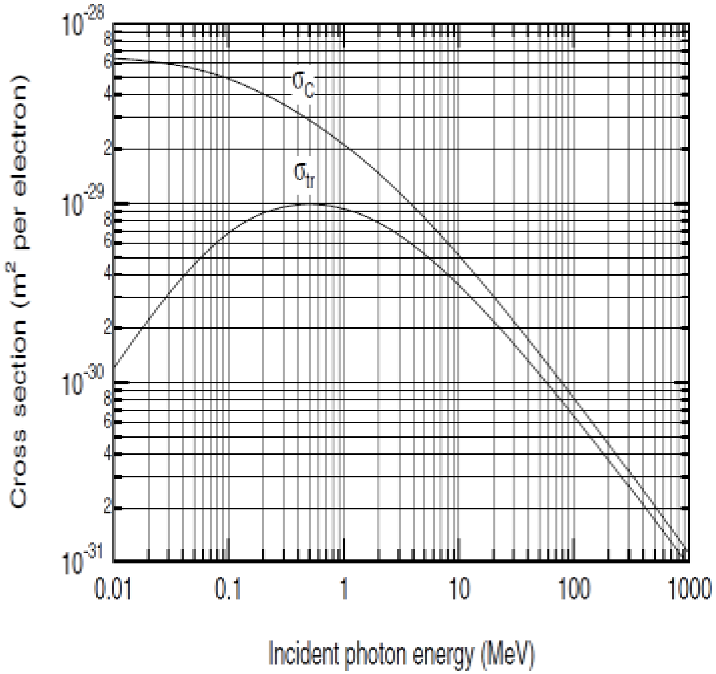
- For carbon, equality near 10 keV.

Compton Scattering: Energy Transferred to Electron

- Integrating T equation over all angles

$$\sigma_{\mathrm{tr}} = \int_0^{\pi} \frac{d\sigma_C}{d\Omega} \frac{T(\theta)}{h\nu_0} 2\pi \sin\theta \, d\theta = f_C \sigma_C .$$

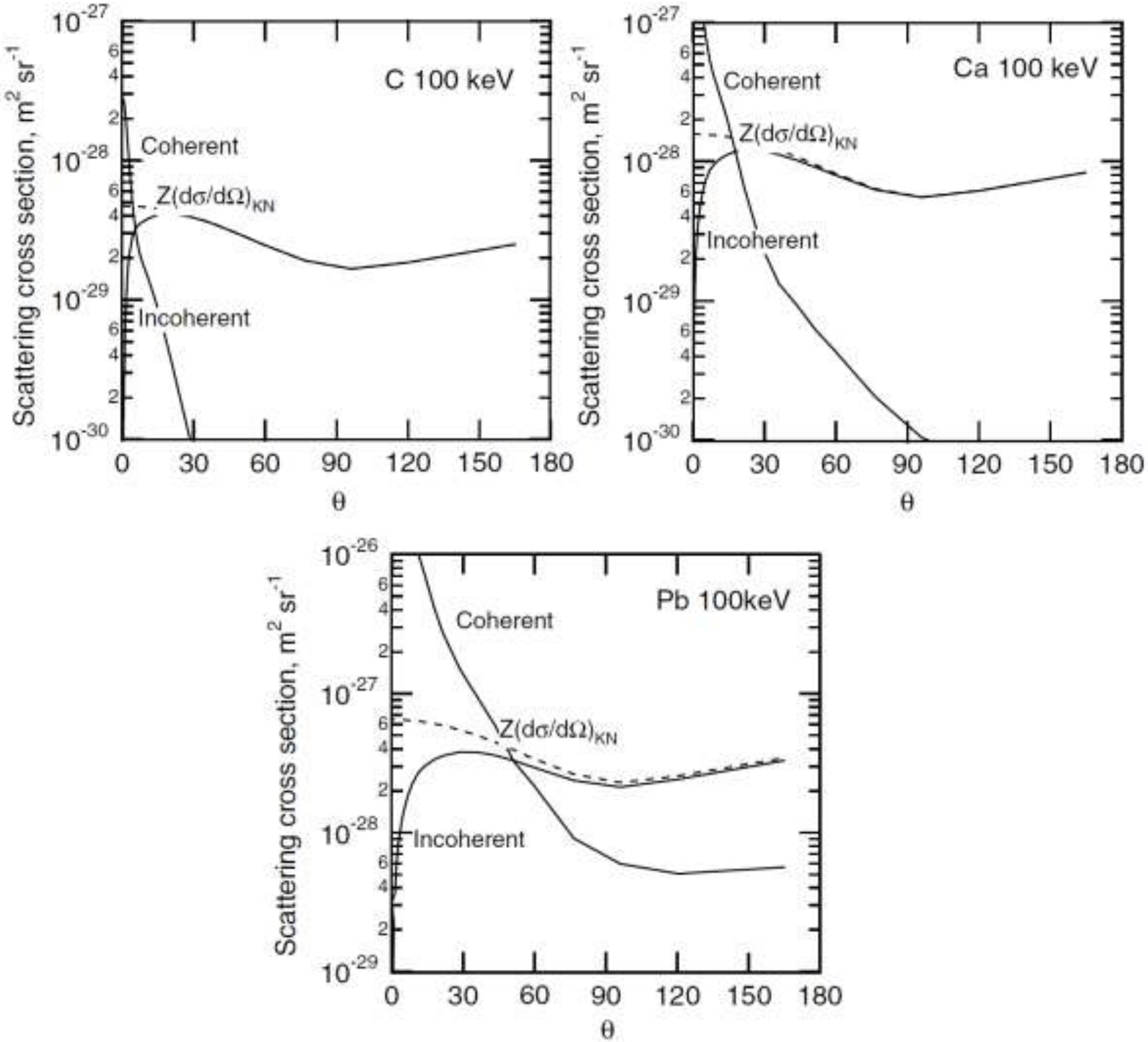
$$\sigma_{\mathrm{tr}} = 2\pi r_e^2 \left[\frac{2(1+x)^2}{x^2(1+2x)} - \frac{1+3x}{(1+2x)^2} - \frac{(1+x)(2x^2-2x-1)}{x^2(1+2x)^2} - \frac{4x^2}{3(1+2x)^3} - \left(\frac{1+x}{x^3} - \frac{1}{2x} + \frac{1}{2x^3} \right) \ln(1+2x) \right].$$



Coherent Scattering

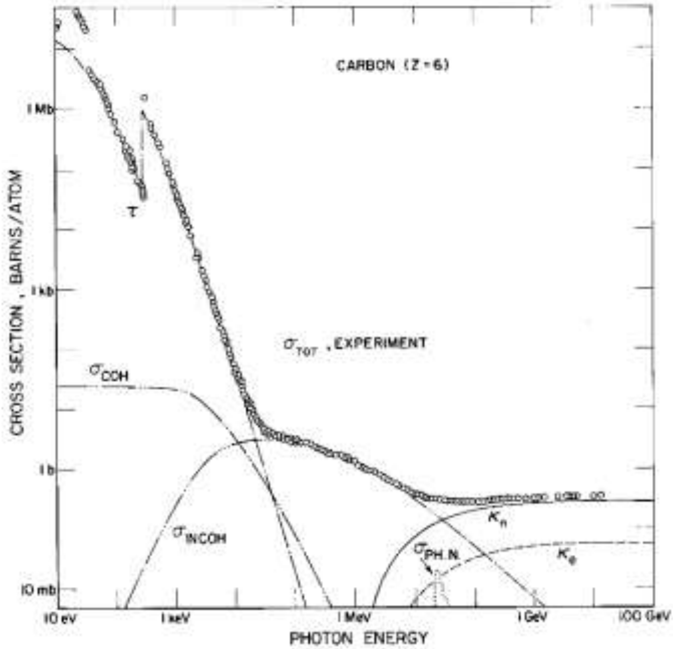
- (γ, γ) photon interaction.
- Primary mechanism is oscillation of electron cloud in the atom in response to the electric field of the incident photons.
- Cross section for coherent scattering is σ_{coh}.
  - σ_{coh} peaked in the forward direction because of interference effects between EM waves scattered by various parts of the electron cloud.
  - Peak is narrower for elements of lower atomic number and for higher energies.

Coherent Scattering



Coherent Scattering

- If wavelength of incident photon >> size of the atom, all Z electrons behave like a single particle with charge −Ze and mass Zm_e.
  - Limit is almost $Z^2\sigma_c$



Pair Production

- High energy (γ,*e*$^+$ *e*$^-$*) interaction*

$$h\nu_0 = T_+ + m_e c^2 + T_- + m_e c^2 = T_+ + T_- + 2m_e c^2.$$

- One can show that momentum is not conserved by the positron and electron if the former equation is satisfied.
  - Interaction takes place in the Coulomb field of another particle (usually a nucleus) that recoils to conserve momentum.
  - Cross section for pair production involving nucleus is κ_n.

Pair Production

- Pair production with excitation or ionization of the recoil atom can take place at energies that are only slightly higher than the threshold
  - Cross section does not become appreciable until the incident photon energy exceeds *2.04 MeV*
  - A free electron (rather than a nucleus) recoils to conserve momentum.
  - (γ*,e*$^+$ *e*$^-$ *e*$^-$*)* process : Triplet production.
- Total cross section: $\kappa = \kappa_n + \kappa_e$

Linear Attenuation Coefficient

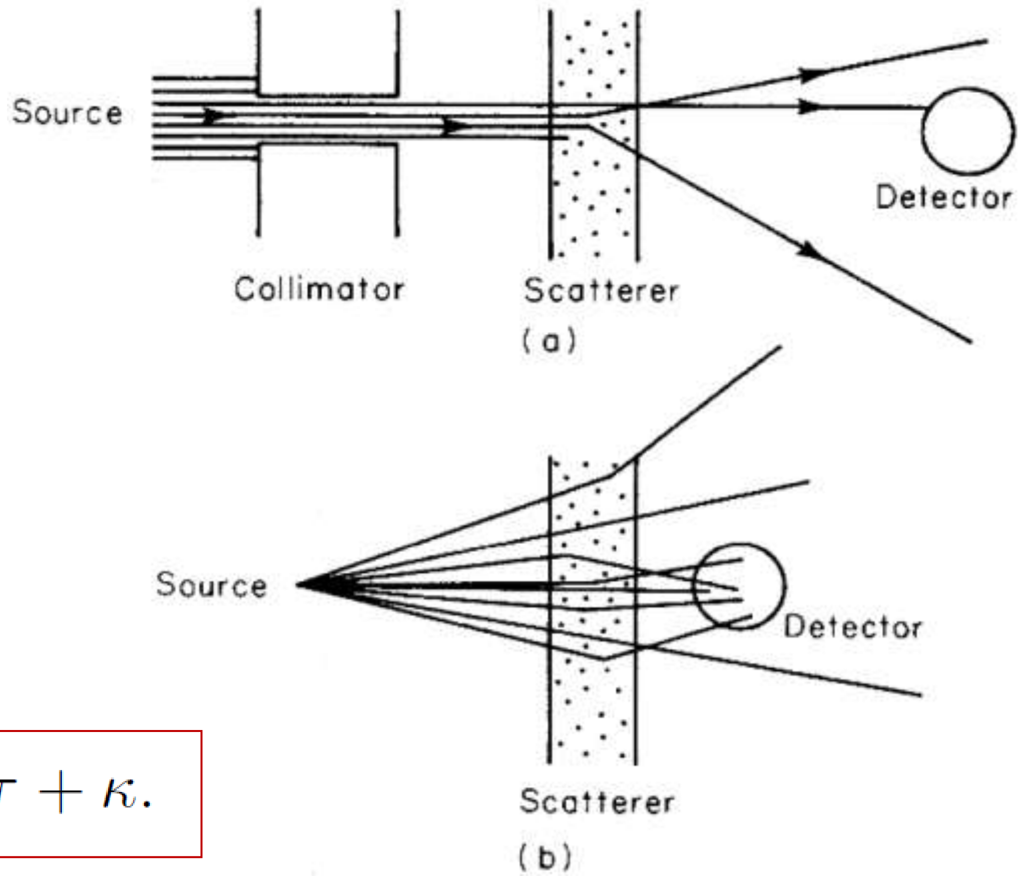
- Narrow- vs. Broad-beam geometries
  - Idealization ?

$$dN = -\frac{\sigma_{\text{tot}} N_A \rho}{A} N\,dz,$$

$$N(z) = N_0 e^{-\mu_{\text{atten}} z}.$$

$$\mu_{\text{atten}} = \frac{N_A \rho \sigma_{\text{tot}}}{A}.$$

$$\sigma_{\text{tot}} = \sigma_{\text{coh}} + \sigma_{\text{incoh}} + \tau + \kappa.$$



Mass Attenuation Coefficient

- Mass attenuation coefficient
  - Independent of density: very useful in gases

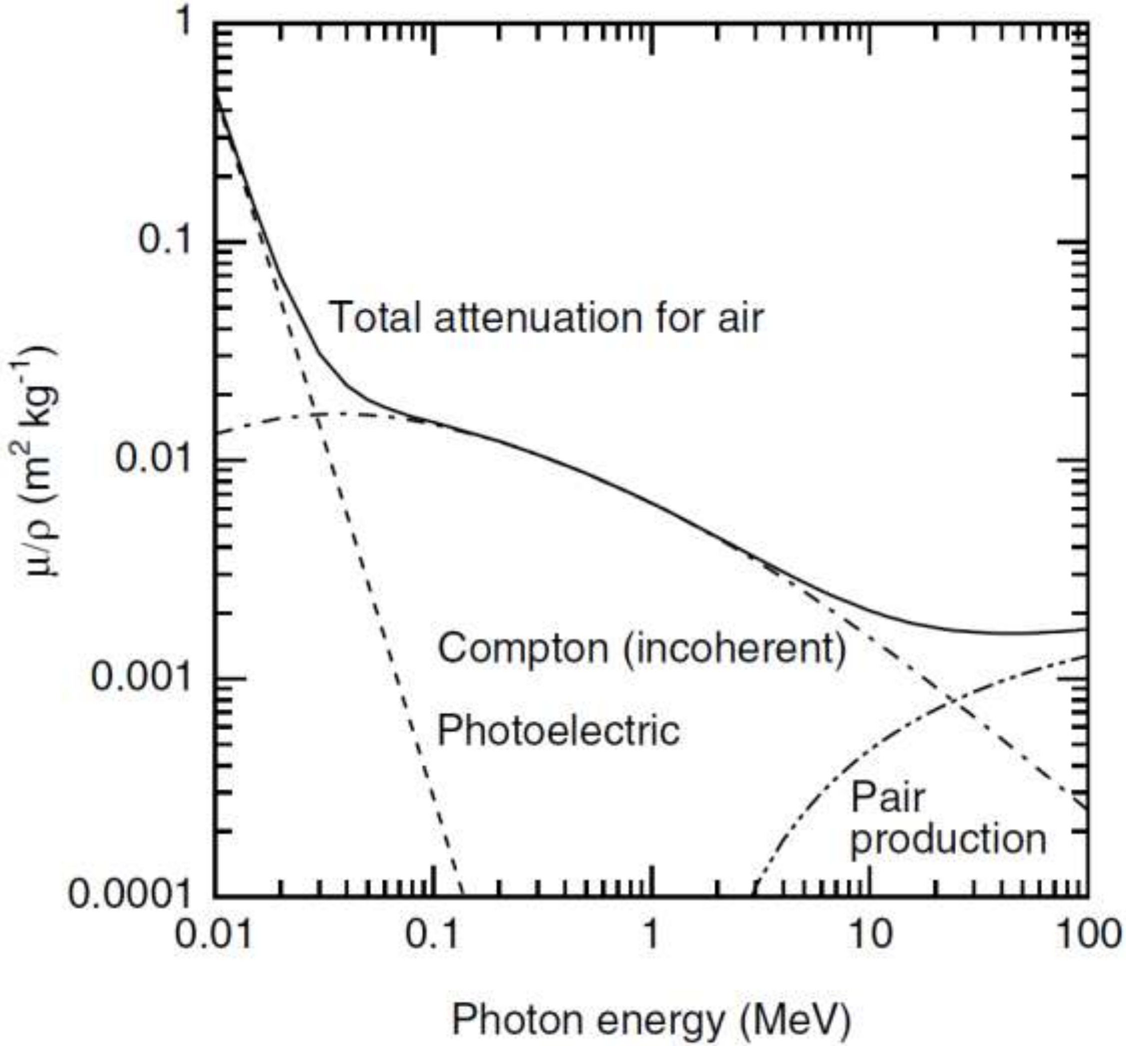
$$\frac{\mu_{\mathrm{atten}}}{\rho} = \frac{N_A \sigma_{\mathrm{tot}}}{A}.$$

$$N(\rho z) = N_0 e^{-(\mu_{\mathrm{atten}}/\rho)(\rho z)}.$$

  - Additional advantage in incoherent scattering: Z/A is nearly ½ for all elements except H^1: minor variations over periodic table

$$\frac{\mu_{\mathrm{atten}}}{\rho} = \frac{Z \sigma_C N_A}{A}$$

Mass Attenuation Coefficient



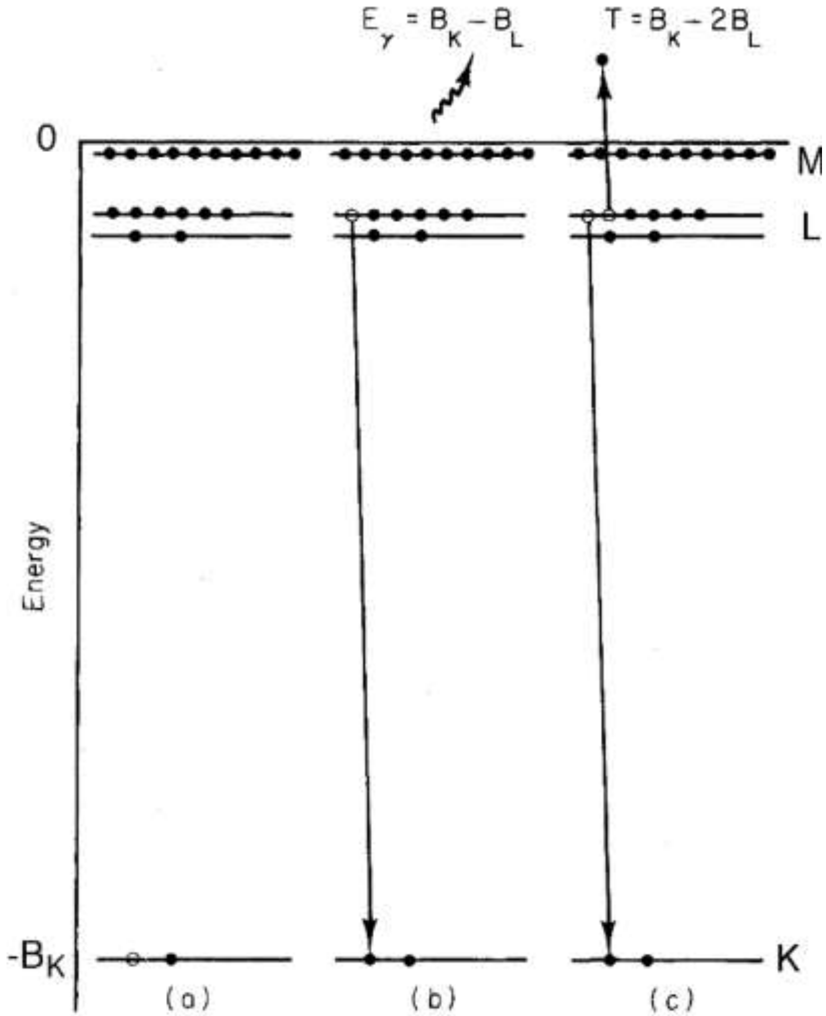
Deexcitation of Atoms

- Excited atom is left with a hole in some electron shell.
  - Similar state when an electron is knocked out by a passing charged particle or by certain transformations in the atomic nucleus
- Two competing processes:
  - Radiative transition: photon is emitted as an electron falls into the hole from a higher level,
  - Nonradiative or radiationless transition: emission of an Auger electron

Deexcitation of Atoms

Process	Total photon energy	Total electron energy	Atom excitation energy	Sum
Before photon strikes atom	$h\nu$	0	0	$h\nu$
After photoelectron is ejected [Fig. 15.12(a)]	0	$h\nu - B_K$	B_K	$h\nu$
Case 1: Deexcitation by the emission of a K and an L photon				
Emission of K fluorescence photon [Fig. 15.12(b)]	$B_K - B_L$	$h\nu - B_K$	B_L	$h\nu$
Emission of L fluorescence photon	$B_K - B_L$, B_L	$h\nu - B_K$	0	$h\nu$
Case 2: Deexcitation by emission of an Auger electron from the L shell				
Emission of Auger electron [Fig. 15.12(c)]	0	$h\nu - B_K$, $B_K - 2B_L$	$2B_L$	$h\nu$
First L-shell hole filled by fluorescence	B_L	$h\nu - B_K$, $B_K - 2B_L$	B_L	$h\nu$
Second L-shell hole filled by fluorescence	B_L, B_L	$h\nu - B_K$, $B_K - 2B_L$	0	$h\nu$

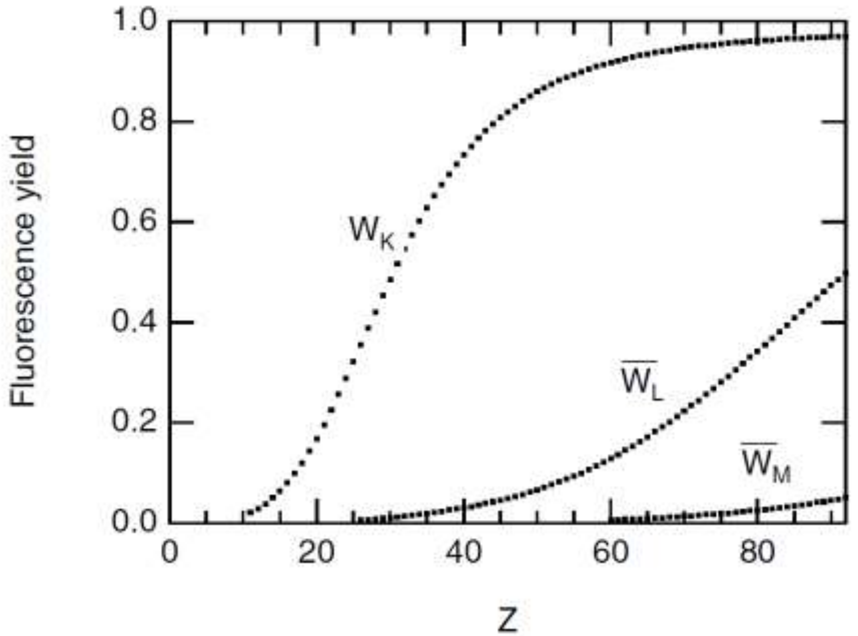
Deexcitation of Atoms



$$\Delta l = \pm 1, \qquad \Delta j = 0, \pm 1.$$

Deexcitation of Atoms

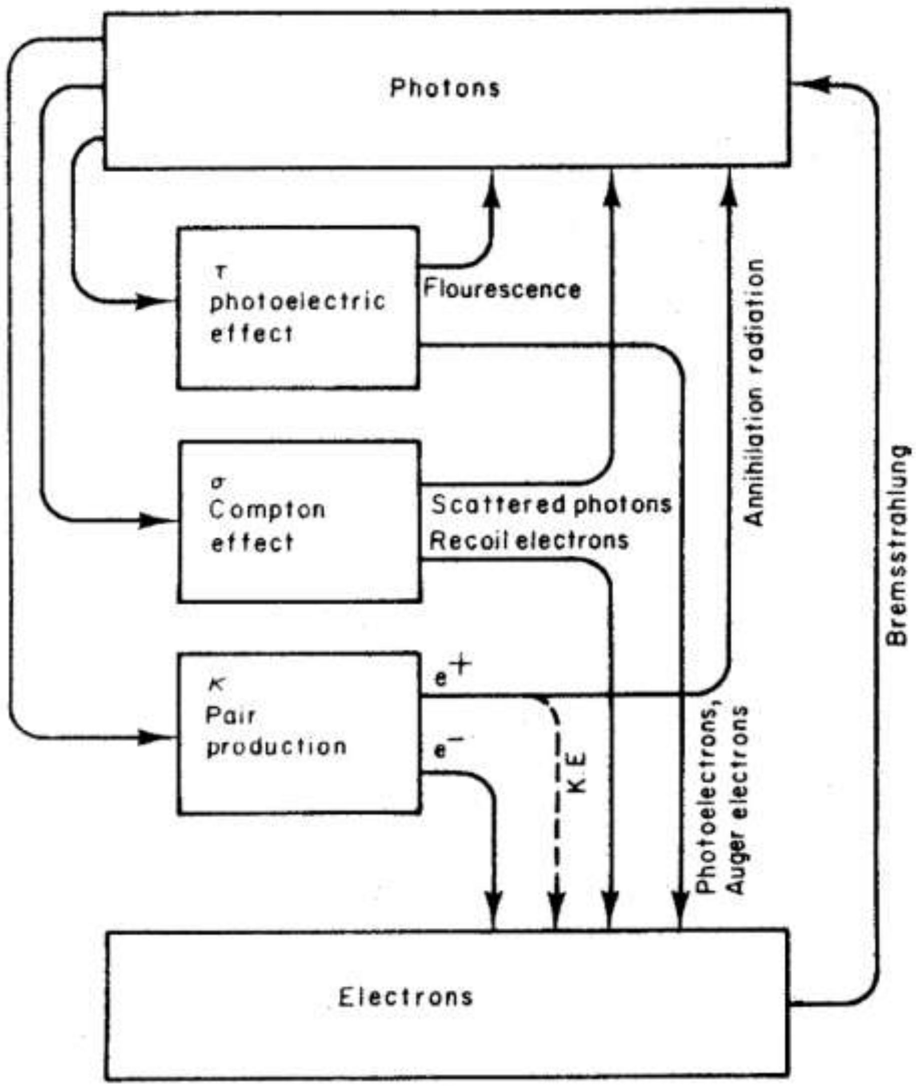
- Probability of photon emission is called the fluorescence yield, W_K.
  - *Auger yield is $A_K = 1 - W_K$.*
  - *L* or higher shells: consider yield for each subshell



Deexcitation of Atoms

- Coster–Kronig transitions
  - Radiationless transitions within the subshell
  - Hole in L_I-shell can be filled by an electron from the L_{III}-shell with the ejection of an M-shell electron
- Super-Coster–Kronig transitions
  - Involves electrons all within same shell (e.g., all M)
- Auger cascade
  - Bond breaking – important for radioactive isotopes

Energy Transfer from Photons to Electrons



Problem Assignments

- Information posted on web site
- Chapter 15 problems: 3, 4, 7, 8, 9, 14, 16, 17, 18, 19, 21